PAPER • OPEN ACCESS

Searching for New Physics in two-neutrino double beta decay with CUPID

To cite this article: A Armatol et al 2021 J. Phys.: Conf. Ser. 2156 012233

View the article online for updates and enhancements.

You may also like

- LOCALIZATION AND BROADBAND FOLLOW-UP OF THE GRAVITATIONAL-WAVE TRANSIENT GW150914 B. P. Abbott, R. Abbott, T. D. Abbott et al.
- A new calibration method for charm jet identification validated with proton-proton collision events at *s* = 13 TeV. The CMS collaboration, Armen Tumasyan, Wolfgang Adam et al.
- Optimization of a single module of CUPID Alberto Ressa, On behalf of the CUPID collaboration, A. Armatol et al.



ECS Membership = Connection

ECS membership connects you to the electrochemical community:

- Facilitate your research and discovery through ECS meetings which convene scientists from around the world;
- Access professional support through your lifetime career:
- Open up mentorship opportunities across the stages of your career;
- Build relationships that nurture partnership, teamwork—and success!

Join ECS!

Visit electrochem.org/join



2156 (2022) 012233 doi:10.1088/1742-6596/2156/1/012233

Searching for New Physics in two-neutrino double beta decay with CUPID

```
A Armatol<sup>1</sup>, E Armengaud<sup>1</sup>, W Armstrong<sup>2</sup>, C Augier<sup>3</sup>,
F T Avignone III<sup>4</sup>, O Azzolini<sup>5</sup>, A Barabash<sup>6</sup>, G Bari<sup>7</sup>, A Barresi<sup>8,9</sup>,
D Baudin<sup>1</sup>, F Bellini<sup>10,11</sup>, G Benato<sup>12</sup>, M Beretta<sup>13</sup>, L Bergé<sup>14</sup>,
M Biassoni<sup>8</sup>, J Billard<sup>3</sup>, V Boldrini<sup>7,15</sup>, A Branca<sup>8,9</sup>, C Brofferio<sup>8,9</sup>,
C Bucci<sup>12</sup>, J Camilleri<sup>16</sup>, S Capelli<sup>8,9</sup>, L Cappelli<sup>12</sup>, L Cardani<sup>10</sup>,
P Carniti<sup>8,9</sup>, N Casali<sup>10</sup>, A Cazes<sup>3</sup>, E Celi<sup>12,17</sup>, C Chang<sup>2</sup>,
M Chapellier<sup>18</sup>, A Charrier<sup>19</sup>, D Chiesa<sup>8,9</sup>, M Clemenza<sup>8,9</sup>,
I Colantoni<sup>10,20</sup>, F Collamati<sup>10</sup>, S Copello<sup>21,22</sup>, O Cremonesi<sup>8</sup>,
R J Creswick<sup>4</sup>, A Cruciani<sup>10</sup>, A D'Addabbo<sup>12,17</sup>, G D'Imperio<sup>10</sup>, I Dafinei<sup>10</sup>, F A Danevich<sup>23</sup>, M de Combarieu<sup>19</sup>, M De Jesus<sup>3</sup>,
P de Marcillac<sup>14</sup>, S Dell'Oro<sup>8,9,16</sup>, S Di Domizio<sup>21,22</sup>, V Dompè<sup>10,11</sup>,
A Drobizhev<sup>24</sup>, L Dumoulin<sup>18</sup>, G Fantini<sup>10,11</sup>, M Faverzani<sup>8,9</sup>,
E Ferri<sup>8,9</sup>, F Ferri<sup>1</sup>, F Ferroni<sup>10,17</sup>, E Figueroa-Feliciano<sup>25</sup>,
J Formaggio<sup>26</sup>, A Franceschi<sup>27</sup>, C Fu<sup>28</sup>, S Fu<sup>28</sup>, B K Fujikawa<sup>24</sup>,
J Gascon<sup>3</sup>, S Ghislandi<sup>12,17</sup>, A Giachero<sup>8,9</sup>, L Gironi<sup>8,9</sup>, A Giuliani<sup>14</sup>, P Gorla<sup>12</sup>, C Gotti<sup>8</sup>, P Gras<sup>1</sup>, M Gros<sup>1</sup>, T D Gutierrez<sup>29</sup>, K Han<sup>30</sup>,
E V Hansen<sup>13</sup>, K M Heeger<sup>31</sup>, D L Helis<sup>1</sup>, H Z Huang<sup>28,32</sup>,
R G Huang<sup>13,24</sup>, L Imbert<sup>18</sup>, J Johnston<sup>26</sup>, A Juillard<sup>3</sup>, G Karapetrov<sup>33</sup>, G Keppel<sup>5</sup>, H Khalife<sup>14</sup>, V V Kobychev<sup>23</sup>,
J Kotila<sup>31,44</sup>, Yu G Kolomensky<sup>13,24</sup>, S Konovalov<sup>6</sup>, Y Liu<sup>34</sup>,
P Loaiza<sup>14</sup>, L Ma<sup>28</sup>, M Madhukuttan<sup>18</sup>, F Mancarella<sup>7,15</sup>, R Mariam<sup>14</sup>,
L Marini<sup>12,13,24</sup>, S Marnieros<sup>14</sup>, M Martinez<sup>35,36</sup>, R H Maruyama<sup>31</sup>,
B Mauri<sup>1</sup>, D Mayer<sup>26</sup>, Y Mei<sup>24</sup>, S Milana<sup>10</sup>, D Misiak<sup>3</sup>,
T Napolitano<sup>27</sup>, M Nastasi<sup>8,9</sup>, X F Navick<sup>1</sup>, J Nikkel<sup>31</sup>, R Nipoti<sup>7,15</sup>,
S Nisi<sup>12</sup>, C Nones<sup>1</sup>, E B Norman<sup>13</sup>, V Novosad<sup>2</sup>, I Nutini<sup>8,9</sup>,
T O'Donnell<sup>16</sup>, E Olivieri<sup>14</sup>, C Oriol<sup>14</sup>, J L Ouellet<sup>26</sup>, S Pagan<sup>31</sup>,
C Pagliarone<sup>12</sup>, L Pagnanini<sup>12,17</sup>, P Pari<sup>19</sup>, L Pattavina<sup>12,37</sup>, B Paul<sup>1</sup>
M Pavan<sup>8,9</sup>, H Peng<sup>38</sup>, G Pessina<sup>8</sup>, V Pettinacci<sup>10</sup>, C Pira<sup>5</sup>, S Pirro<sup>12</sup>,
D V Poda<sup>14</sup>, T Polakovic<sup>2</sup>, O G Polischuk<sup>23</sup>, S Pozzi<sup>8,9</sup>, E Previtali<sup>8,9</sup>,
A Puiu<sup>12,17</sup>, S Quitadamo<sup>12,17</sup>, A Ressa<sup>10,11</sup>, R Rizzoli<sup>7,15</sup>, C Rosenfeld<sup>4</sup>,
C Rusconi<sup>12</sup>, V Sanglard<sup>3</sup>, J Scarpaci<sup>14</sup>, B Schmidt<sup>24,25</sup>, V Sharma<sup>16</sup>,
V Shlegel<sup>39</sup>, V Singh<sup>13</sup>, M Sisti<sup>8</sup>, D Speller<sup>31,40</sup>, P T Surukuchi<sup>31</sup>,
L Taffarello<sup>41</sup>, O Tellier<sup>1</sup>, C Tomei<sup>10</sup>, V I Tretyak<sup>23</sup>, A Tsymbaliuk<sup>5</sup>,
M Velazquez<sup>42</sup>, K J Vetter<sup>13</sup>, S L Wagaarachchi<sup>13</sup>, G Wang<sup>2</sup>,
L Wang<sup>34</sup>, B Welliver<sup>24</sup>, J Wilson<sup>4</sup>, K Wilson<sup>4</sup>, L A Winslow<sup>26</sup>,
M Xue<sup>38</sup>, L Yan<sup>28</sup>, J Yang<sup>38</sup>, V Yefremenko<sup>2</sup>, V Yumatov<sup>6</sup>,
M M Zarytskyy<sup>23</sup>, J Zhang<sup>2</sup>, A Zolotarova<sup>14</sup>, S Zucchelli<sup>7,43</sup>
```

¹IRFU, CEA, Université Paris-Saclay, Saclay, France

²Argonne National Laboratory, Argonne, IL, USA

Content from this work may be used under the terms of the Creative Commons Attribution 3.0 licence. Any further distribution of this work must maintain attribution to the author(s) and the title of the work, journal citation and DOI.

2156 (2022) 012233 doi:10.1088/1742-6596/2156/1/012233

```
<sup>3</sup>Institut de Physique des 2 Infinis, Lyon, France
```

E-mail: emanuela.celi@gssi.it

Abstract. In the past few years, attention has been drawn to the fact that a precision analysis of two-neutrino double beta decay $(2\nu\beta\beta)$ allows the study of interesting physics cases like the emission of Majoron bosons and possible Lorentz symmetry violation. These processes modify the summed-energy distribution of the two electrons emitted in $2\nu\beta\beta$. CUPID is a next-generation experiment aiming to exploit 100 Mo-enriched scintillating Li₂MoO₄ crystals, operating as cryogenic calorimeters. Given the relatively fast half-life of 100 Mo $2\nu\beta\beta$ and the large exposure that can be reached by CUPID, we expect to measure with very high precision the 100 Mo $2\nu\beta\beta$ spectrum shape, reaching great sensitivities in the search for distortions induced by the physics beyond the Standard Model. In this contribution, we present the CUPID exclusion sensitivity for such New Physics processes, as well as the preliminary projected background of CUPID.

⁴University of South Carolina, Columbia, SC, USA

⁵INFN Laboratori Nazionali di Legnaro, Legnaro, Italy

⁶National Research Centre Kurchatov Institute, Institute for Theoretical and Experimental Physics, Moscow, Russia

⁷INFN Sezione di Bologna, Bologna, Italy

⁸INFN Sezione di Milano - Bicocca, Milan, Italy

⁹University of Milano - Bicocca, Milan, Italy

¹⁰INFN Sezione di Roma, Rome, Italy

¹¹Sapienza University of Rome, Rome, Italy

¹²INFN Laboratori Nazionali del Gran Sasso, Assergi (AQ), Italy

¹³University of California, Berkeley, CA, USA

¹⁴Université Paris-Saclay, CNRS/IN2P3, IJCLab, Orsay, France

¹⁵CNR-Institute for Microelectronics and Microsystems, Bologna, Italy

¹⁶Virginia Polytechnic Institute and State University, Blacksburg, VA, USA

¹⁷Gran Sasso Science Institute, L'Aquila, Italy

¹⁸Université Paris-Saclay, CNRS/IN2P3, IJCLab, Orsay, France

¹⁹IRAMIS, CEA, Université Paris-Saclay, Saclay, France

²⁰CNR-Institute of Nanotechnology, Rome, Italy

²¹INFN Sezione di Genova, Genova, Italy

 $^{^{22} \}mathrm{University}$ of Genova, Genova, Italy

²³Institute for Nuclear Research of NASU, Kyiv, Ukraine

 $^{^{24} \}mathrm{Lawrence}$ Berkeley National Laboratory, Berkeley, CA, USA

²⁵Northwestern University, Evanston, IL, USA

²⁶Massachusetts Institute of Technology, Cambridge, MA, USA

²⁷INFN Laboratori Nazionali di Frascati, Frascati, Italy

²⁸Fudan University, Shanghai, China

²⁹California Polytechnic State University, San Luis Obispo, CA, USA

³⁰Shanghai Jiao Tong University, Shanghai, China

³¹Yale University, New Haven, CT, USA

³²University of California, Los Angeles, CA, USA

 $^{^{33}\}mathrm{Drexel}$ University, Philadelphia, PA, USA

 $^{^{34} \}mbox{Beijing Normal University, Beijing, China}$

³⁵Universidad de Zaragoza, Zaragoza, Spain

³⁶ARAID Fundación Agencia Aragonesa para la Investigación y el Desarrollo, Zaragoza, Spain

³⁷Physik-Department, Technische Universität München, Garching, Germany

³⁸University of Science and Technology of China, Hefei, China

³⁹Nikolaev Institute of Inorganic Chemistry, Novosibirsk, Russia

⁴⁰Johns Hopkins University, Baltimore, MD, USA

⁴¹INFN Sezione di Padova, Padova, Italy

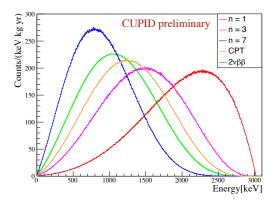
⁴²Univ. Grenoble Alpes, CNRS, Grenoble INP, SIMAP, Grenoble, France

⁴³University of Bologna, Bologna, Italy

⁴⁴University of Jyväskylä, Jyväskylä, Finland

Journal of Physics: Conference Series

2156 (2022) 012233 doi:10.1088/1742-6596/2156/1/012233



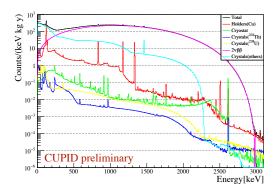


Figure 1: Standard $2\nu\beta\beta$ decay compared to bSM spectra.

Figure 2: CUPID projection of the β/γ spectrum.

1. Introduction

Although the discovery of neutrinoless double beta decay $(0\nu\beta\beta)$ is the primary target of CUPID [1], other interesting physics cases can be investigated, in particular those beyond the Standard Model (bSM) processes that can induce a deviation of the $2\nu\beta\beta$ spectral shape. Indeed, the phase space factor G depends on the spectral index n through the relation $G \sim (Q_{\beta\beta} - T)^n$ where $Q_{\beta\beta}$ is the Q-value of the decay and T is the summed kinetic energies of the two emitted electrons. For standard $2\nu\text{DBD}$ n=5, while for bSM processes the spectral index assumes different values inducing a shift of the maximum of the spectrum as shown in Fig. 1. An accurate background projection of CUPID is crucial to perform the sensitivity studies on the $2\nu\beta\beta$ spectral shape. In Sec 2 the preliminary CUPID Background Budget is shown in details, while in Sec 3 we explain the analysis method used to evaluate the CUPID exclusion sensitivity for the CPT violating $2\nu\beta\beta$ and for several Majoron emitting decays.

2. The CUPID Background Budget

The Background Budget (BB) is composed by a series of Monte Carlo simulations aiming to predict the CUPID background spectrum. We simulate radioactive contaminations using a Geant4 [2] based software. Our knowledge about radioactive contaminants in CUPID comes mainly from material assays, previous bolometric experiments and cosmogenic activation calculations. Since CUPID will be hosted in the same cryogenic infrastructure of CUORE, the CUORE background budget [3] provides a clear picture of contaminations in the detector holders and the cryostat, while the CUPID-Mo data describe the impurities of Li_2MoO_4 crystals [4]. Combining these models, we have the possibility to assess almost all the background sources, having a reliable estimate of the CUPID background in a wide energy region. Besides, we introduced in the BB also contaminations due to cosmogenic activation in the Li_2MoO_4 crystals and copper holders that were calculated with the ACTIVIA code [5]. The simulations are then processed with a custom software to implement experimental features on simulated data, like the energy and the time resolution, the coincidence window and the particle identification. The projection of the CUPID β/γ spectrum is shown in Fig. 2.

3. Exclusion sensitivity results

Using the same software tools, we simulate the energy spectrum of bSM processes starting from the exact phase space calculation for $2\nu\beta\beta$ [6], obtaining the results reported in Fig. 1. To evaluate the CUPID sensitivity for a given bSM process after 1 yr of data-taking, we simulate the corresponding statistics according to the BB. Performing a combined Bayesian fit on CUPID-like

Journal of Physics: Conference Series

2156 (2022) 012233

doi:10.1088/1742-6596/2156/1/012233

Table 1: List of decays with one (χ^0) and two $(\chi^0\chi^0)$ Majoron emission.

\overline{n}	mode	exclusion sensitivity on $T_{1/2}$ [yr]	current limit [yr] (NEMO-3)
1	$\chi_{_{0}}^{0}$	7.4×10^{23}	$4.4 \times 10^{22} [8]$
3	χ^0	2.4×10^{21}	4.4×10^{21} [9]
3	$\chi^0 \chi^0 $	$\begin{array}{c} 2.4 \times 10^{21} \\ 7.3 \times 10^{21} \end{array}$	4.4×10^{21} [9]
7	χ ° χ °	7.3×10^{21}	$1.2 \times 10^{21} [9]$

data including the New Physics hypotheses, we set limits on the half-life of bSM processes. The fitting procedure is accurately described in Ref. [7]. One of the critical points in this fit is represented by the pure β -decaying isotopes which can correlate with $2\nu\beta\beta$ and bSM spectra in the fit. 90 Sr is an anthropogenic radioactive isotope that decays with a half-life of 28.8 yr in 90 Sr \rightarrow 90 Y \rightarrow 90 Zr emitting two subsequently β -decays with Q-values, respectively, of 546 keV and 2.3 MeV. The preliminary CUPID-Mo background model assessed that the 90 Sr activity in Li₂MoO₄ crystals is $\sim 10^{-4}$ Bq/kg, but the presence of 90 Sr in the actual CUPID crystals is not certain. To estimate the effect of the pure β -decays on the sensitivity, the analysis was repeated in two cases: the 90 Sr is included in the simulated spectrum but not considered in the fit (underestimation) and, on the contrary, the 90 Sr is considered in the fit but not included in the simulated spectrum (overestimation).

Several grand unification theories predict that one or two Majorons could be emitted in the $0\nu\beta\beta$ [10] producing a continuum spectrum similar to $2\nu\beta\beta$ but with different spectral indexes. The preliminary exclusion sensitivities are shown in Tab. 1. We report the less stringent exclusion sensitivity obtained for each process in the ⁹⁰Sr overestimated and underestimated fits.

The Standard Model Extension (SME) provides a general framework for Lorentz Invariance Violation (LIV) [11]. The parameter $\mathbf{\mathring{a}}_{of}^{(3)}$ is related to the time-like component of the LIV operator in the neutrino sector. The preliminary predicted exclusion limit of the Lorentz-violating term is $\mathbf{\mathring{a}}_{of}^{(3)} \lesssim 10^{-8}$ GeV at 90% C.I., while the current limit was set by NEMO-3 $\mathbf{\mathring{a}}_{of}^{(3)} < 3.5 \times 10^{-7}$ GeV at 90% CL. [9].

References

- [1] Armstrong W R et al. (CUPID) 2019 (Preprint 1907.09376)
- [2] Agostinelli S et al. (GEANT4) 2003 Nucl. Instrum. Meth. A 506 250–303
- [3] Alduino C et al. (CUORE) 2017 Eur. Phys. J. C 77 543 (Preprint 1704.08970)
- [4] Armengaud E et al. 2020 Eur. Phys. J. C 80 674 (Preprint 1912.07272)
- [5] Back J J and Ramachers Y A 2008 Nucl. Instrum. Meth. A **586** 286–294 (Preprint 0709.3472)
- [6] Kotila J and Iachello F 2012 Phys. Rev. C 85 034316 (Preprint 1209.5722)
- [7] Alduino C et al. (CUORE) 2017 Eur. Phys. J. C 77 13 (Preprint 1609.01666)
- [8] Arnold R et al. (NEMO-3) 2015 Phys. Rev. D 92 072011 (Preprint 1506.05825)
- [9] Arnold R et al. (NEMO-3) 2019 Eur. Phys. J. C 79 440 (Preprint 1903.08084)
- [10] Georgi H M, Glashow S L and Nussinov S 1981 Nucl. Phys. B 193 297–316
- [11] Kostelecky V A and Samuel S 1989 Phys. Rev. D 39 683